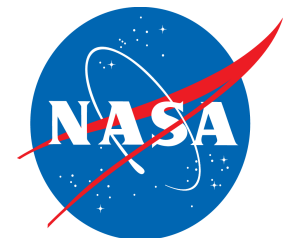


Application of a Physical Continuum Model to Recent X-ray Observations of Accreting Pulsars

Diana M. Marcu-Cheatham

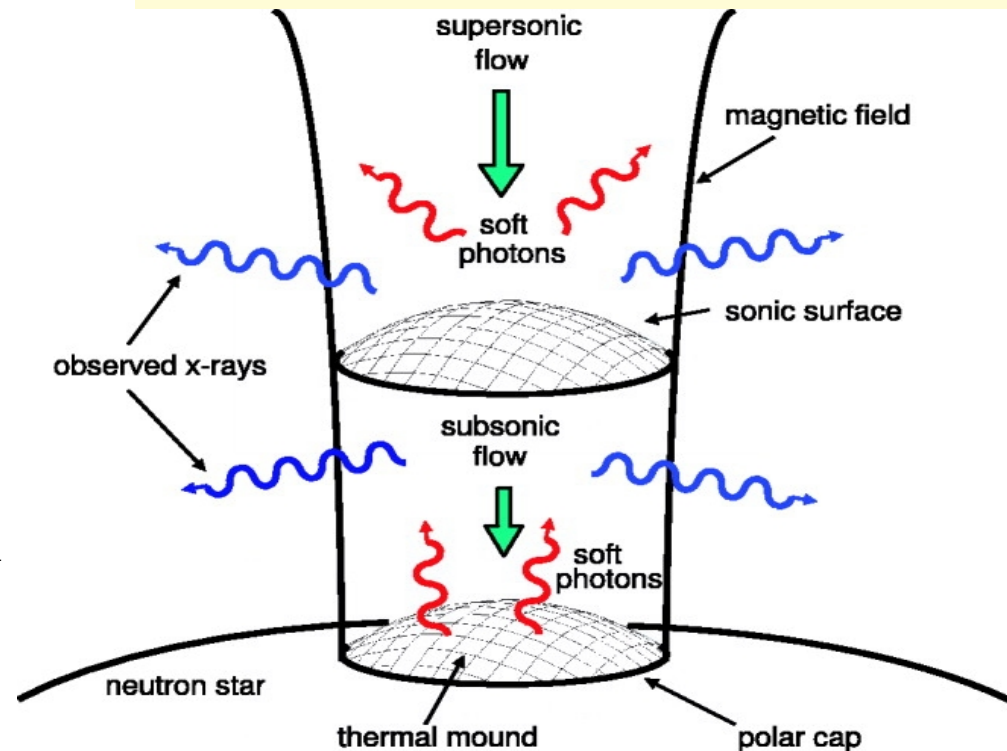
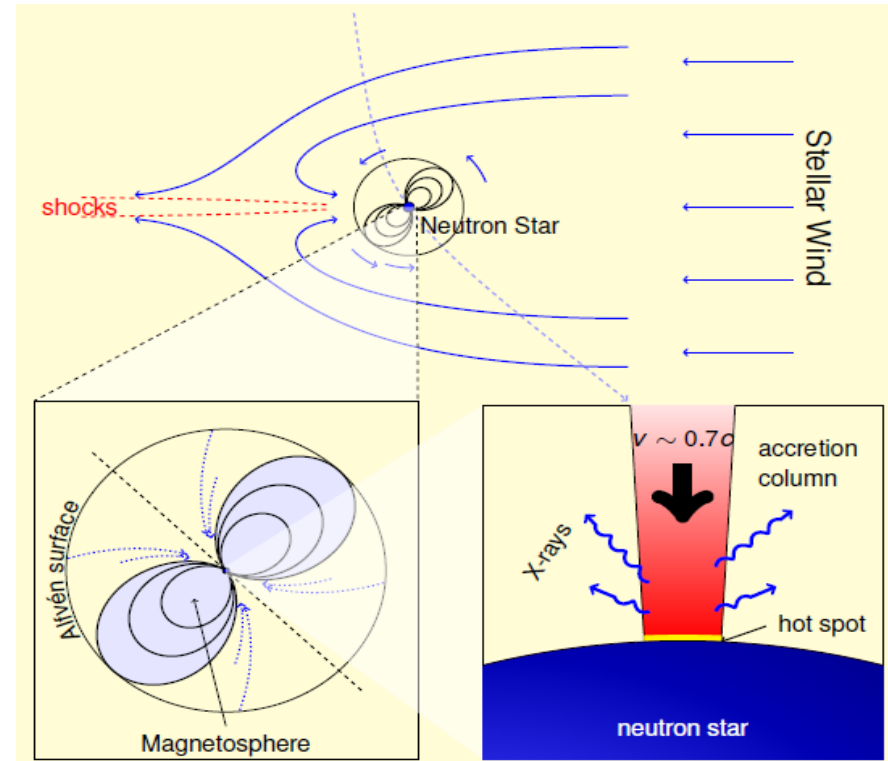
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Wood, J. Wilms, P.B. Hemphill, A.M. Gottlieb,
F. Fuerst, F.-W. Schwarm, R. Ballhausen

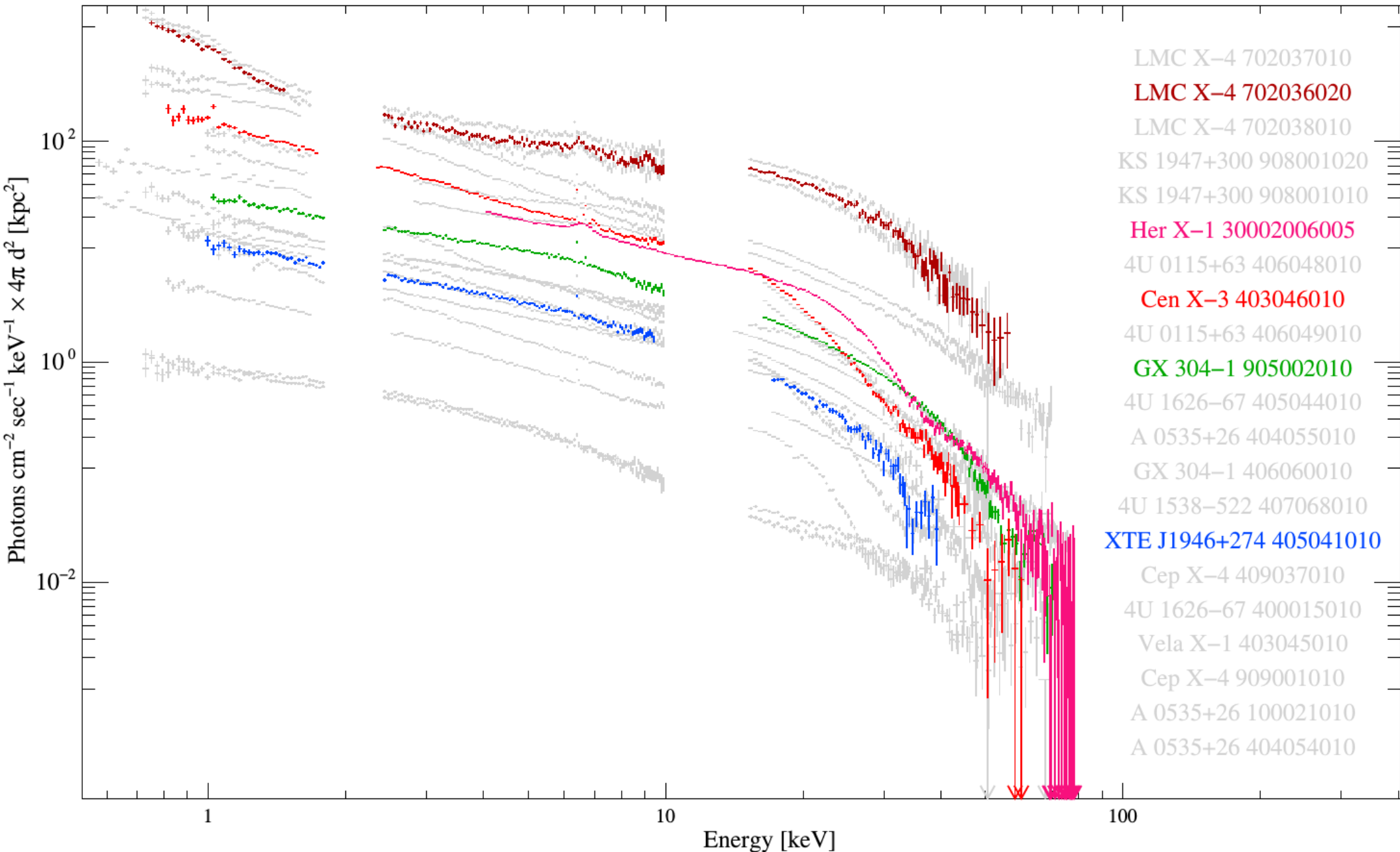


Accretion onto NS Surface

- X-ray seed photons emission in the accretion column:
 - Bremsstrahlung emission
 - Cyclotron emission
 - Blackbody emission
- Seed photons are **up-scattered by electron scattering** in the accretion column:
 - **Bulk Comptonization** due to the bulk motion of infalling electrons
 - **Thermal Comptonization** due to stochastic electron motion
 - Solution to **Radiative Transfer Equation** using **Analytical Formalism** by **Becker & Wolff 2007**
- Empirical models are used to describe cutoff power law shape of the continuum (e.g. Cutoff Power Law, Fermi Dirac Cutoff, etc.)



Sources and Properties

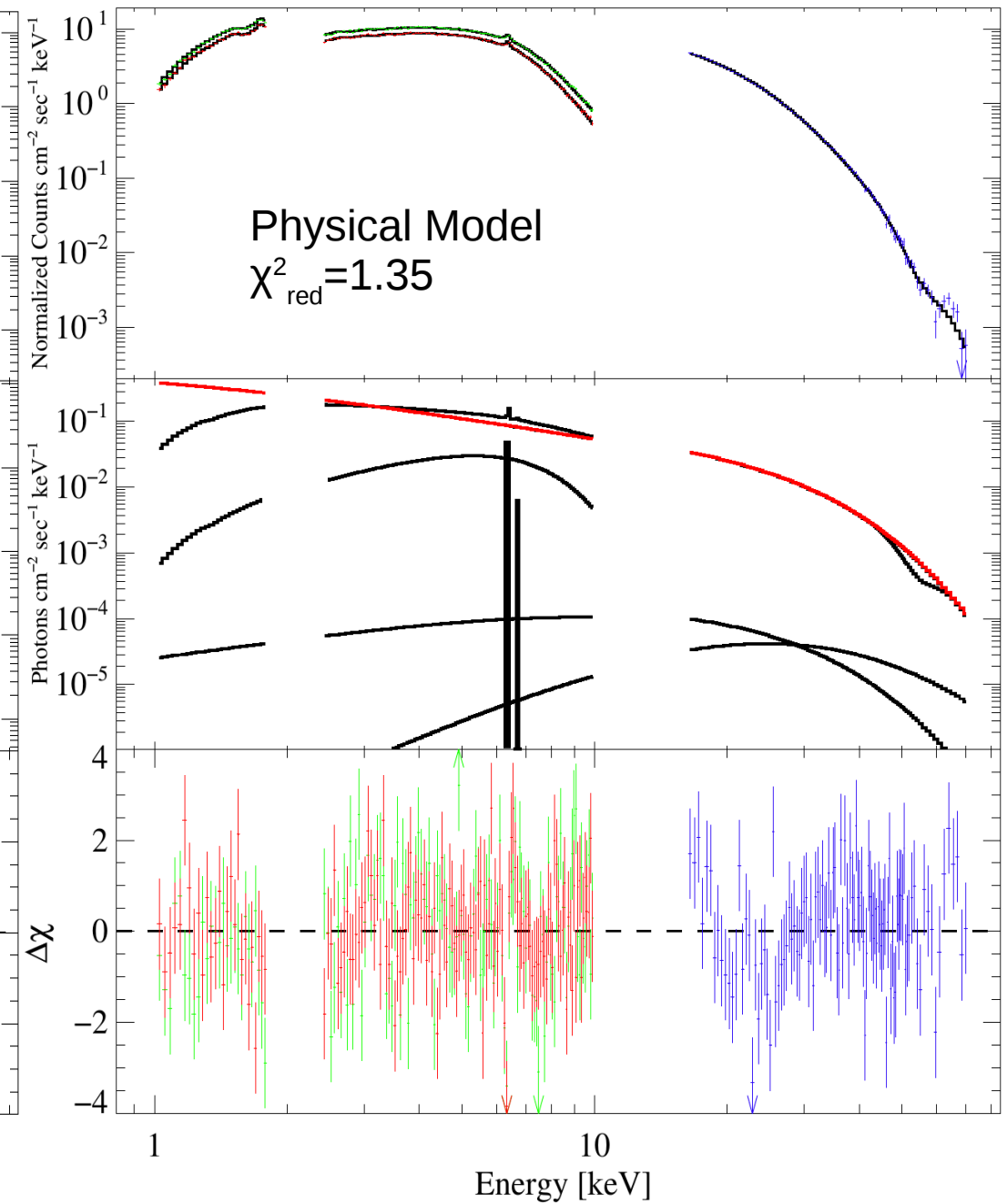
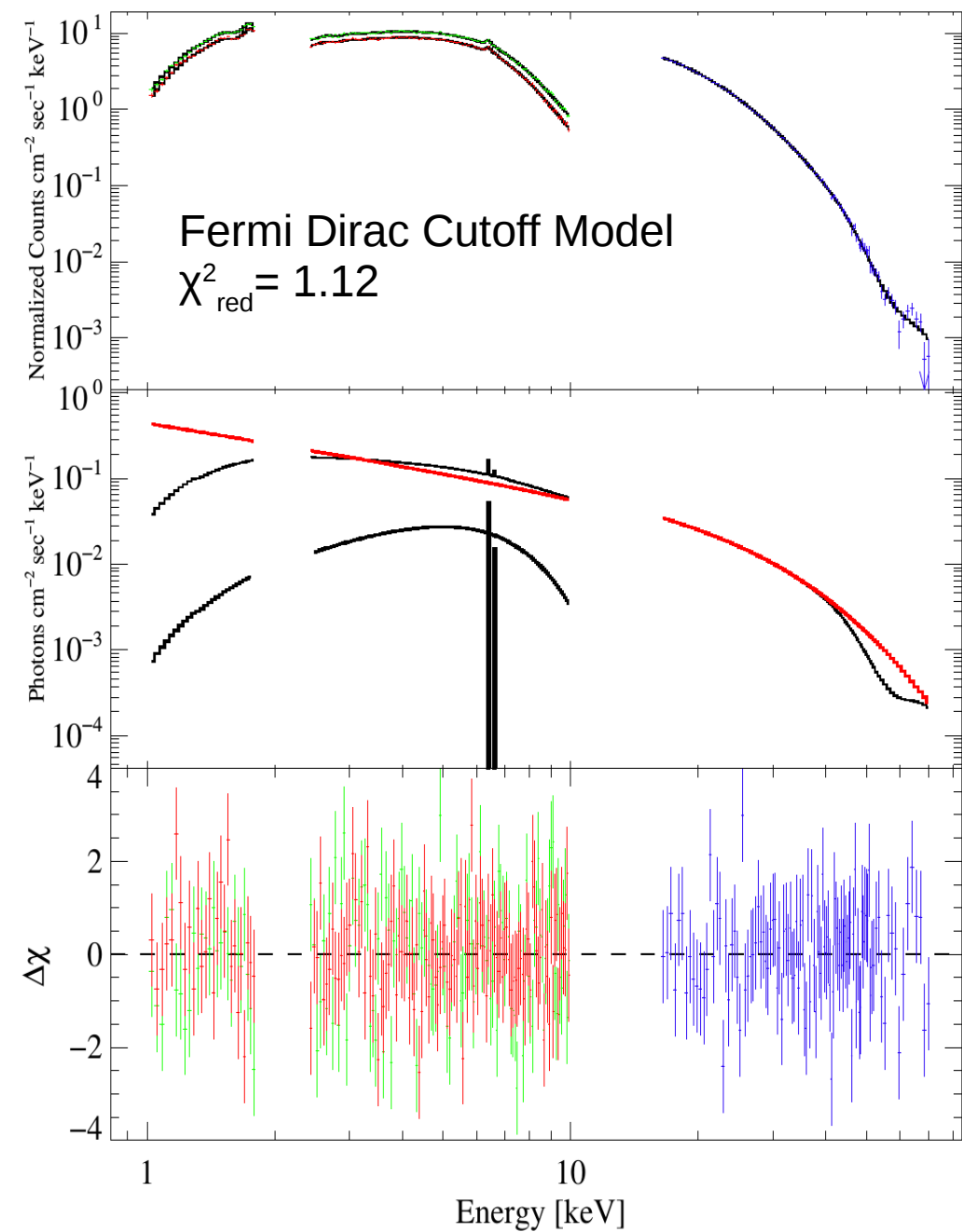


- GX 304-1
- XTE J1946+274

- Cen X-3 (Gottlieb et al. 2016 in prep.) Poster #120.09
- Her X-1 (Wolff et al. 2016 Submitted) Poster #120.24

- All sources have measured cyclotron lines
- except LMC X-4

Spectral Fitting GX 304-1



Fit results

Physical & Fermi Dirac Cutoff Models

Source L_X 1–80keV	kT_e [keV]	r_0 [m]	σ_{\parallel} [$\times 10^{-5} \sigma_T$]	$\bar{\sigma}$ [$\times 10^{-4} \sigma_T$]	E_{fold} [keV]	E_{cut} [keV]	Γ
LMC X-4 $25.4 \times 10^{37} \text{ erg s}^{-1}$	$6.53^{+0.08}_{-0.15}$	1912^{+759}_{-604}	$1.14^{+1.12}_{-0.52}$	106.06	$8.0^{+0.8}_{-0.7}$	$26.2^{+2.0}_{-2.7}$	0.65(4)
1 Her X-1 ^{a,b} $4.9 \times 10^{37} \text{ erg s}^{-1}$	$4.58(7)^a$	$170.4^{+1.7}_{-1.8}{}^a$	$5.2(1)^a$	$3.5(2)^a$	$6.0^{+0.5}_{-0.8}{}^b$	$35.9^{+3.0}_{-0.3}{}^b$	$0.888^{+0.008}_{-0.009}{}^b$
Cen X-3 ^c $4.0 \times 10^{37} \text{ erg s}^{-1}$	$3.1^{+0.4}_{-0.1}{}^c$	$65^{+12}_{-4}{}^c$	$0.28(2)^c$	$1.6^{+0.6}_{-0.3}{}^c$	$7.0^{+0.7}_{-0.6}{}^c$	18(3) ^c	$1.1^{+0.2}_{-0.1}{}^c$
GX 304–1 $2.0 \times 10^{37} \text{ erg s}^{-1}$	6.28	69.4	2.7	8.5	$9.3^{+0.8}_{-0.7}$	$29.2^{+2.4}_{-2.7}$	1.47(3)
XTE J1946+274 $0.5 \times 10^{37} \text{ erg s}^{-1}$	4.45(3)	14.0(2)	1.07(2)	$7.0^{+0.5}_{-0.4}$	$9.1(4)^d$	0.0^d	$0.59(2)^d$

Note. — (a) Wolff et al. 2016 submitted, Poster [#120.24](#); (b) Fuerst et al. (2013); (c) Gottlieb et al. 2016 in prep, Poster [#120.09](#); (d) Marcu-Cheatham et al. (2015).

Conclusions

- Results:
 - All sources have statistically good fits $1.04 < \chi^2_{\text{red}} < 1.35$
 - Regardless of the diversity of the sources, there are **no strong variations in T_e : $\sim 3 - 6$ keV** – similarly to **E_{fold} : $\sim 6 - 9$ keV**
 - Possible correlation: $r_0 \sim L_x$
 - Physical fits reproduce the empirical (FDCO) fits \rightarrow features that strongly influence the continuum were found in both, i.e., BBODY (GX 304-1), “10 keV bump” (Cen X-3), strong Fe line (Her X-1)
- Revolutionizing spectral modeling of neutron star accretion
 - Large improvement in modeling the X-ray continuum by successfully **statistically fitting a physical model**
 - Estimates for accretion column parameters: T_e , r_0 , cross sections
- Next: Physics of Cyclotron Lines

Her X-1, Cen X-3, LMC X-4

- Comparison with BW 2007:
 - T_e approximately the same in both cases, even though luminosities are different
 - Largest differences are in r_0 , especially in LMC X-4, other two show smaller differences probably due to different luminosities
 - Cen X-3 data was taken at a much higher luminosity in BW 2007

Comparison with Farinelli et al. 2015 (F15)

- F15 states that their velocity profile is more generalized rather than the free fall velocity used by BW
- F15 calculates the second order (vs. first order in BW) bulk-Comptonization → T_e slightly lower in F15
- F15 claims a high cyclotron emission that described the “10 keV bump”, however, we don't see a comparably strong cyclotron emission
- F15 finds much higher r_0 values (on the order of 1 – 3 km)

Empirical Continuum Models

$$\text{CUTOFFPL}(E) = AE^{-\Gamma} \times e^{-E/E_{\text{fold}}}$$

$$\text{FDCO}(E) = AE^{-\Gamma} \frac{1}{1 + e^{(E-E_{\text{cut}})/E_{\text{fold}}}}$$

$$\text{PLCUT}(E) = AE^{-\Gamma} \times \begin{cases} 1 & (E \leq E_{\text{cut}}) \\ e^{-(E-E_{\text{cut}})/E_{\text{fold}}} & (E > E_{\text{cut}}) \end{cases}$$

$$\text{NPEX}(E) = (A_1 E^{-\alpha_1} + A_2 E^{+\alpha_2}) e^{-E/E_{\text{fold}}}$$

Empirical CRSF Models

$$M(E) = \begin{cases} \text{GABS}(E) = \tau_c e^{-(E-E_c)^2/(2\sigma_c^2)} \\ \text{CYCLABS}(E) = D_c \frac{(W_c E/E_c)^2}{(E-E_c)^2 + W_c^2} \end{cases}$$

$$I_0(E) \rightarrow I_0(E) e^{-M(E)}$$

Becker and Wolff Physical Model for Accreting Pulsars

- Green's function distribution f_G satisfies the generalized Kompaneets (1957) transport equation

$$\begin{aligned} v \frac{\partial f_G}{\partial z} &= \frac{dv}{dz} \frac{\epsilon}{3} \frac{\partial f_G}{\partial \epsilon} + \frac{\partial}{\partial z} \left(\frac{c}{3n_e \sigma_{\parallel}} \frac{\partial f_G}{\partial z} \right) - \frac{f_G}{t_{\text{esc}}} \\ &+ \frac{n_e \bar{\sigma} c}{m_e c^2} \frac{1}{\epsilon^2} \frac{\partial}{\partial \epsilon} \left[\epsilon^4 \left(f_G + kT_e \frac{\partial f_G}{\partial \epsilon} \right) \right] + \frac{\dot{N}_0 \delta(\epsilon - \epsilon_0) \delta(z - z_0)}{\pi r_0^2 \epsilon_0^2} \end{aligned}$$

- f_G convolved with a source distribution function for each type of emission provides the final spectral model

Becker and Wolff Physical Model for Accreting Pulsars

- Green's function distribution f_G satisfies the generalized Kompaneets (1957) transport equation

$$\begin{aligned}
 \text{Advection} \quad v \frac{\partial f_G}{\partial z} &= \text{First order Fermi energization "bulk" Comptonization} \quad \frac{dv}{dz} \frac{\epsilon}{3} \frac{\partial f_G}{\partial \epsilon} + \text{Spatial diffusion along column axis} \quad \frac{\partial}{\partial z} \left(\frac{c}{3n_e \sigma_{\parallel}} \frac{\partial f_G}{\partial z} \right) - \text{Photon escape} \quad \frac{f_G}{t_{\text{esc}}} \\
 &+ \text{Thermal Comptonization} \quad \frac{n_e \bar{\sigma} c}{m_e c^2} \frac{1}{\epsilon^2} \frac{\partial}{\partial \epsilon} \left[\epsilon^4 \left(f_G + kT_e \frac{\partial f_G}{\partial \epsilon} \right) \right] + \text{Photon injection} \quad \frac{\dot{N}_0 \delta(\epsilon - \epsilon_0) \delta(z - z_0)}{\pi r_0^2 \epsilon_0^2}
 \end{aligned}$$

NOTE! Final model is LINEAR: $F_{\epsilon}(\epsilon) \equiv \frac{[\Phi_{\epsilon}^{\text{cyc}}(\epsilon) + \Phi_{\epsilon}^{\text{bb}}(\epsilon) + \Phi_{\epsilon}^{\text{ff}}(\epsilon)] A_c(\epsilon)}{4\pi D^2}$